

## 1. Problem set 2: Integrability

Problem 1:

For the Sinh-Gordon theory, solve the asymptotic Bethe Ansatz equation in the two particle sector, for the case of zero total momentum. What is the energy of the lowest order solution ? How does it depend on  $m$  and  $L$  ? For  $mL \gg 1$ , do they behave like fermions? Why? What is the low energy scattering phase?. Can you compute the correction to the energy for an excited state with relatively large momentum ?. The  $S$  matrix is

$$S(\theta) = \frac{\sinh \theta - i \sin \pi p}{\sinh \theta + i \sin \pi p}, \quad 0 < p \leq 1/2 \quad (1.1)$$

Does this  $S$  matrix have a bound state?. Compare to the scattering of similar massive particles in Sinh Gordon which gives something like (1.1) but with the opposite sign for  $p$ . For the sine gordon can you find the energy of the bound states ?.

Problem 2:

Show that a single massive scalar field in two dimensions with interaction  $\phi^5$  is not integrable. ( Hint: does the  $2 \rightarrow 3$  amplitude vanish ? )

Problem 3:

Consider a chain of spins. (Spin 1/2 representations of  $SU(2)$ ). Consider the most general  $SU(2)$  invariant nearest neighbor Hamiltonian. Show that, up to a constant shift, it is the Heisenberg one. Show that is is given by the permutation operator. Why is this?. Hint use group theory and the decomposition of two spin 1/2 representations into spin one and zero.

## 2. $\mathcal{N} = 4$ super Yang Mills spin chains

Problem 1:

Consider the operator  $\dots ZZZZWZZZZ \dots$  with an  $W$  with some momentum. Understand the computation of the one loop anomalous dimension. Consider the diagrams that come from the term in the potential of the form  $Tr\{[Z, W][\bar{Z}, \bar{W}]\}$ . Understand the two terms that appear and their relative sign. Then compute the logarithmic terms in the correction to the two point function of operators  $\sum_l e^{ip_l} Tr[.ZZZZWZZZ\dots]$ , where  $W$  is at position  $l$ , and its complex conjugate? Here we have been vague about the  $\dots$ . You can consider a concrete operator such as  $\sum_l Tr[WZ^l WZ^{J-l}]e^{ip_l}$  and compute the anomalous dimension.

Problem 2:

Look up in Wess and Bagger the supersymmetry transformations of a theory with a superpotential. Conclude that  $\delta_Q W \sim \psi$  and  $\delta_Q \psi \sim [Z, W]$ .

Problem 4:

Show that in a closed chain the sum of central charges  $\kappa_i$  vanishes, once we take into account the extra phase factors and momentum conservation  $e^{i \sum_{i=1}^n p_i} = 1$ . In other words the full generator is

$$\hat{K} \propto \sum_j f_j (1 - e^{ip_j}) , \quad \text{where} \quad f_i = \prod_{k=1}^{j-1} e^{ip_k} \quad (2.1)$$

where the sum is over all the impurities in the chain.

Imagine that you assume that  $\hat{K}$  has some expression of the form

$$\hat{K} \propto \sum_j f_j h(p_j) , \quad \text{where} \quad f_i = \prod_{k=1}^{j-1} e^{ip_k} \quad (2.2)$$

with  $h$  some unknown function of  $p_j$  (and  $\lambda$ ). The factor of  $f_j$  have to emerge since we know that the  $\{QQ\}$  anticommutator inserts a  $Z$ . Now let's fix  $h$  as follows. We first make the further assumption that for small  $p$   $h(p)$  is linear in  $p$ . (Why is this a very reasonable assumption?). Then we will demand that on any closed chain with a set of particles conserving momentum we should have that  $\hat{K}$  is zero. Choosing appropriate sets of particles show that  $h$  has the form in (2.1). (Hint: Consider a particle of momentum  $p_1$  which is arbitrary,  $0 < p_1 < 2\pi$ . Then consider particles  $p_2, \dots, p_n$  with infinitesimal momentum, but that sums to  $2\pi - p_1$ , and take the large  $n$  limit. Use the fact that  $h$  is linear in  $p$  for small  $p$  and replace the sums by an integrals.)

This exercise shows that the functional form of  $h$  is fixed to (2.1) for all values of the coupling. Have we used integrability in this argument ?

Problem 4:

Consider the  $SU(2|2)$  algebra with the extra central charges. The full algebra is given in (2.5) and (2.6) of hep-th/0511082 . Check that the fundamental representation has the energy stated in the lectures. That is shown in the same paper, but you can check it yourselves.

### 3. References

A great reference on integrability of 1+1 dimensional systems and the construction of the asymptotic S matrices for sine Gordon and the  $O(n)$  models is:

*Factorized s Matrices in Two-Dimensions as the Exact Solutions of Certain Relativistic Quantum Field Models.* A. B. Zamolodchikov and Al. B. Zamolodchikov, Annals Phys.120:253-291,1979.

A nice reference on spin chains and their symmetries is:

*Algebraic aspects of Bethe Ansatz* , L.D. Faddeev, hep-th/9404013

A good place to start reading about the  $N = 4$  Yang Mills spin chain is:

*The  $SU(2|2)$  dynamic S-matrix* , N. Beisert, hep-th/0511082

Table 1 in hep-th/0511082 contains the full matrix form of the  $S$  matrix. The phase factor is given in an integral form in (46) (47) of hep-th/0703104. It was derived originally in hep-th/0609044 and hep-th/0610251

The computation of the cusp anomalous dimension can be found in hep-th/0610251 (but you will need to have read hep-th/0603157 ).